

# G<sub>2</sub>-QCD at Finite Density

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collaboration with

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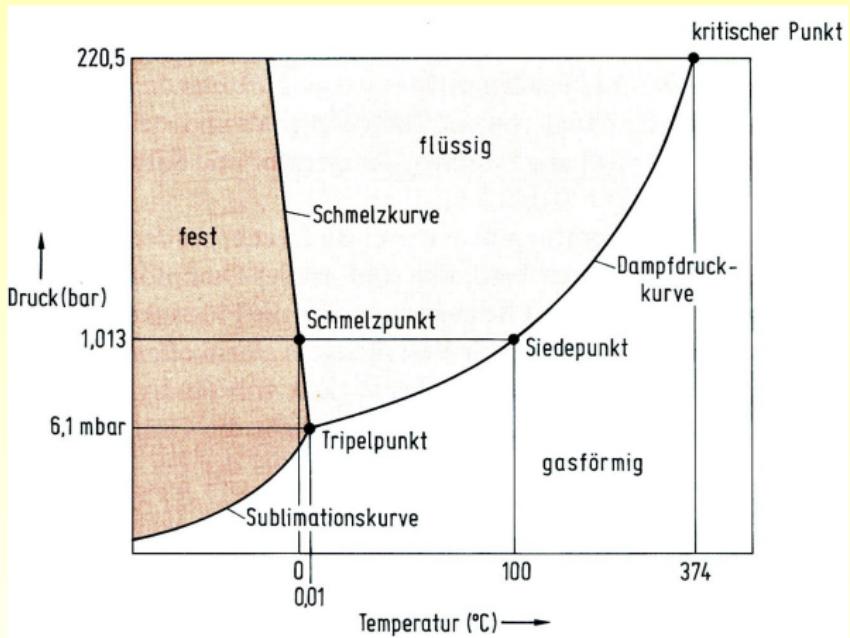
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- 1 Introduction
- 2  $G_2$  and  $G_2$  Gauge Theory
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# Water

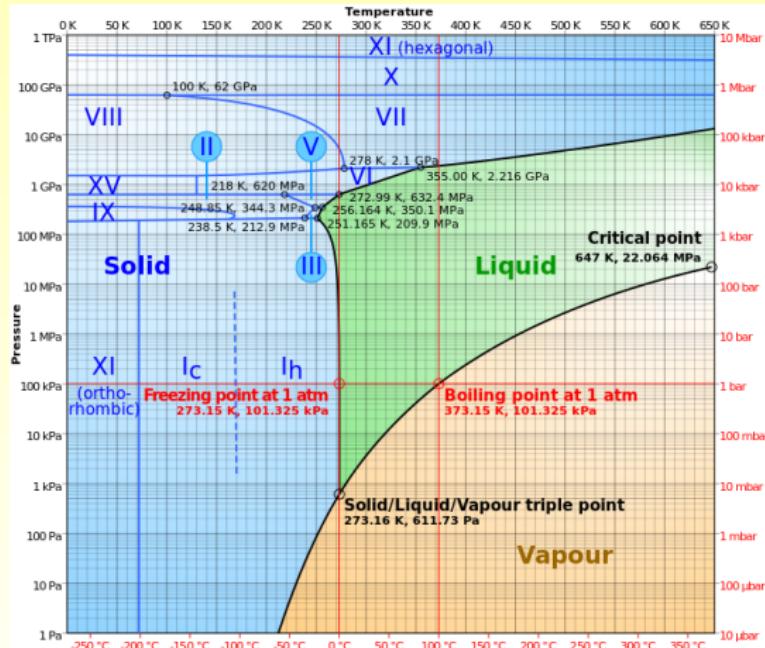


different phases, transition lines, critical point, triple point, . . .



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that is not all:



- crystalline and amorphous phases
- order of H-binding
- critical point
- triple points
- structural transitions
- small  $p$ : hex/cub ice, ice XI
- high  $p$ : ice VII, VIII, X



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# Some facts about exceptional Lie group $G_2$

- smallest exceptional Lie group
- rank = 2, dimension = 14
- real group (not only pseudo-real!)
- subgroup of  $SO(7)$
- $G_2/SU(3) \sim S^7 \rightarrow$  efficient parametrization
- fundamental representations  $\{7\}$ ,  $\{14\}$  (= adjoint)  
7 quarks instead of 3 (cp. GUTS)
- can be broken to  $SU(3)$  with scalars in  $\{7\}$   
fermions:  $\{7\} \rightarrow \{3\} + \{\bar{3}\} + \{1\}$ , gauge bosons:  $\{14\} \rightarrow \{8\} + X$
- smallest (simply connected) Lie group with trivial center



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- singlet representation colorless states

$$\{7\} \otimes \{7\} = \{1\} \oplus \{7\} \oplus \{14\} \oplus \{27\}$$

$$\{7\} \otimes \{7\} \otimes \{7\} = \{1\} \oplus 4 \cdot \{7\} \oplus 2 \cdot \{14\} \oplus \dots$$

$$\{14\} \otimes \{14\} = \{1\} \oplus \{14\} \oplus \{27\} \oplus \dots,$$

$$\{14\} \otimes \{14\} \otimes \{14\} = \{1\} \oplus \{7\} \oplus 5 \cdot \{14\} \oplus \dots,$$

$$\{7\} \otimes \{14\} \otimes \{14\} = \{1\} \oplus \dots$$

- branching  $G_2 \rightarrow SU(3)$

$$\{7\} \longrightarrow \{3\} \oplus \{\bar{3}\} \oplus \{1\},$$

$$\{14\} \longrightarrow \{8\} \oplus \{3\} \oplus \{\bar{3}\}.$$



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# Same facts about $G_2$ gauge theory

- asymptotically free
- first order confinement/deconfinement PT Pepe et al.; Greensite; Cossu et al.
- no center, no order parameter
- chiral restoration in quenched theory at same  $T_c$  Graz group
- qualitatively similar glueball spectrum as SU(3) Welleghausen, Wozar, AW
- Casimir scaling to high accuracy Welleghausen, Wozar, AW; Liptat et al.; Greensite et al.
- topological properties, instantons Maas, Olejnik, Ilgenfritz



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# Why consider G<sub>2</sub> Theories

- G<sub>2</sub> has trivial center:  
confinement models based on center?
- G<sub>2</sub> contains SU(3) as subgroup:  
smooth interpolation between G<sub>2</sub>- and SU(3)-gauge theories
- particle spectrum comparable to QCD (mesons and baryons + ...)
- G<sub>2</sub> has no sign problem:  
simulations at finite  $T$  and finite possible  $\mu$  possible
- G<sub>2</sub>-QCD has baryons and mesons:  
can build a "neutron star"



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- understanding of  $G_2$  under extreme conditions
- phases and phase-transition at finite  $T$  and  $n_B$
- distinguish phases:  
densities, pressure, energy density, condensates, symmetries, ...  
order parameters  $\Leftrightarrow$  symmetries
- vary control parameters:  
temperature, chemical potentials, fields,
- physics question:  
what are the relevant degrees of freedom in a given phase?



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# Lattice simulations for SU(3)

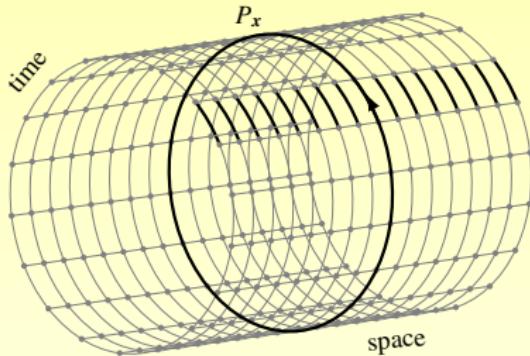
- large baryon density:  
sign problem  
⇒ not accessible to simulations based on important sampling
- effective models, functional methods, ....
- proposals/speculations on exotic phases of cold dense matter
- relevant of  $n^*$ ?
- simulations of theories without sign problem
- even better: solve sign problem?
- here: fast implementation of (local)HMC, no low acceptance rate,



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# Confinement in pure $G_2$ gauge theory

finite  $T \Rightarrow$   
lattice = cylinder with  
circumference  
 $\beta_T = 1/kT = aN_0$ .



- approximate order parameter: Polyakov loop

$$P(x) = \text{tr } \mathcal{P} \exp \left[ i \int_0^\beta d\tau A_0(\tau, x) \right]$$

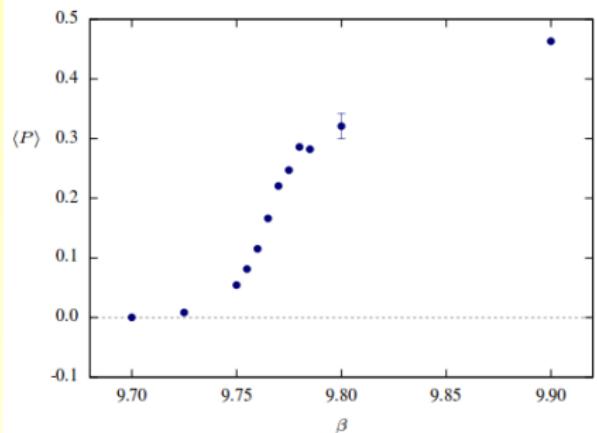
- static potential

$$\langle P(x) \rangle_\beta = e^{-\beta F(x)} , \quad \langle P(x) P^\dagger(y) \rangle_\beta = e^{-\beta V_{q\bar{q}}(R)}$$

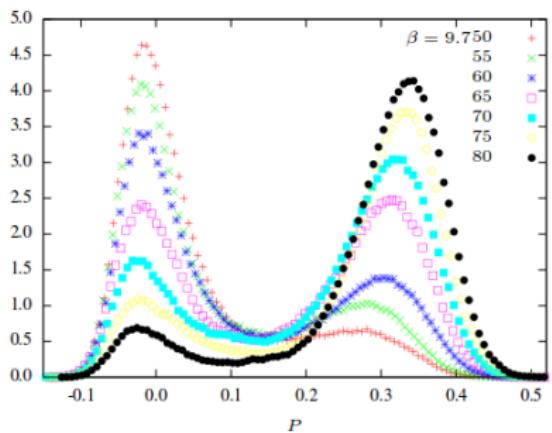


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## Polyakov loop in fundamental representation



rapid change with  $\beta = 1/g^2$



histogramm in vicinity of  $\beta_c$

- Polyakov loop approximate order parameter
- first order PT as in SU(3) gluodynamics



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- $V_{q\bar{q}}(R)$  static potential

confinement:  $\langle P \rangle = 0$  , de-confinement:  $\langle P \rangle \neq 0$

- confinement:  $V \rightarrow \sigma R \Rightarrow \langle P(x)P^\dagger(y) \rangle_\beta \propto e^{-\sigma \cdot \text{Area}}$
- $e^{\sigma \cdot \text{Area}}$  varies over 100 orders of magnitude
- brute force approach does not work
- Lüscher and Weisz method: exponential error reduction
- split lattice in time slices
- calculate  $\langle \dots \rangle$  with fixed bc on each slice
- full result: integral over boundary conditions
- iteration → *multilevel algorithm*



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- static potential for charges in representation  $\mathcal{R}$ :

$$V_{\mathcal{R}}(R) = \gamma_{\mathcal{R}} - \frac{\alpha_{\mathcal{R}}}{R} + \sigma_{\mathcal{R}} R$$

- Casimir scaling hypothesis for string tensions:

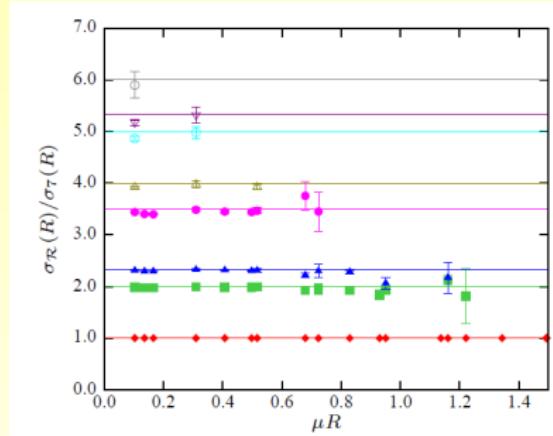
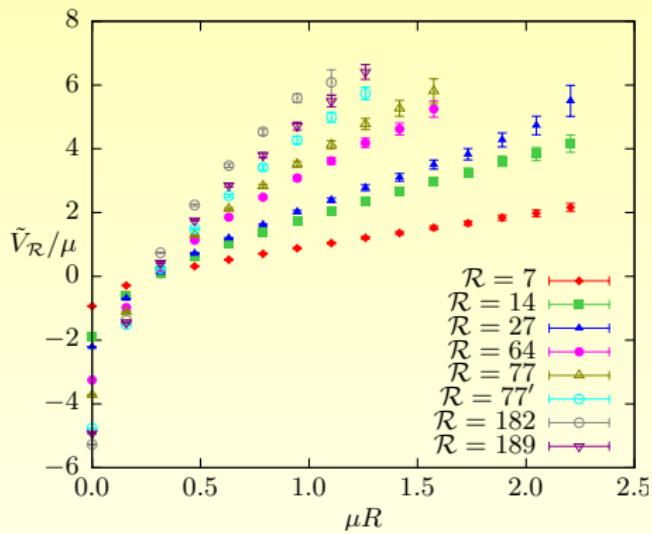
$$\frac{\sigma_{\mathcal{R}}}{c_{\mathcal{R}}} = \frac{\sigma_{\mathcal{R}'}}{c_{\mathcal{R}'}}$$

- from ratios of Wilson(Polyakov) loops

$$V_{\mathcal{R}}(R) = \frac{1}{\tau} \ln \frac{\langle W_{\mathcal{R}}(R, T) \rangle}{\langle W_{\mathcal{R}}(R, T + \tau) \rangle}.$$



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Scaled local string Tension with  $\beta = 9.7, 10$  on  $14^4, 20^4$

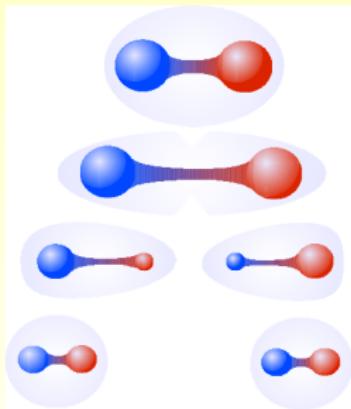
- linear potential for static quarks in different  $G_2$  representations

Welleghausen, AW., Wozar



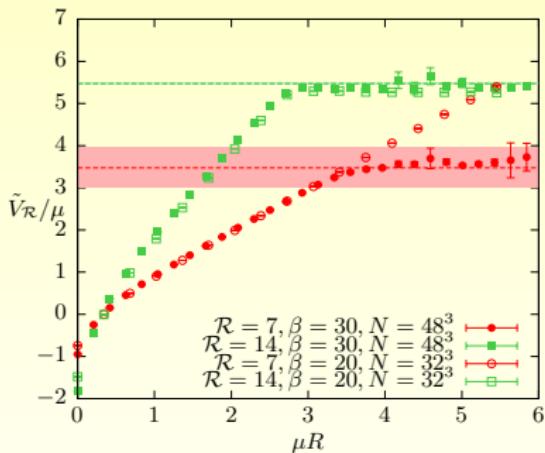
# String-breaking

- meson, diquark  $\bar{q}q \rightarrow$  2 mesons, diquarks



Click here

- energy scale =  $2 m_{\text{glueball}}$
- decay products: glue-lumps



Welleghausen, AW., Wozar (2011)



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# $G_2$ Yang-Mills-Higgs theory

- breaking  $G_2 \rightarrow SU(3)$
- lattice action with normalized Higgs  $\varphi = (\varphi_1, \dots, \varphi)^T$  in  $\{7\}$

$$S_{\text{YMH}}[\mathcal{U}, \varphi] = -\frac{1}{g^2} \sum \mathcal{U}_\square - \kappa \sum \varphi_x^T \mathcal{U}_{x,\mu} \varphi_{x+\mu}$$

- Higgs-mechanism for  $v = \langle \varphi \rangle \neq 0$ :
- $\{14\} \longrightarrow \{8\} \oplus \{3\} \oplus \{\bar{3}\}$ 
  - $\{8\}$ :  $SU(3)$  gluons
  - $\{3\} + \{\bar{3}\}$ : massive Vector bosons
- scalars  $7 \rightarrow 1$



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- $\kappa = 0$ : pure  $G_2$  gauge theory:  
first order deconfinement transition
- $\kappa = \infty$ : 6 vector bosons decouple, pure  $SU(3)$   
first order deconfinement transition
- first order transition line connecting two theories?
- calculate  
Polyakov loop and plaquette actions, susceptibilities  
on grid in  $\beta \propto 1/g^2$ ,  $\kappa$ -plane ( $\beta = 5 \dots 10$ ,  $\kappa = 0 \dots 10^4$ )  $\Rightarrow$

### Phase diagram of $G_2$ YM theory



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# average actions and susceptibilities (small lattice)

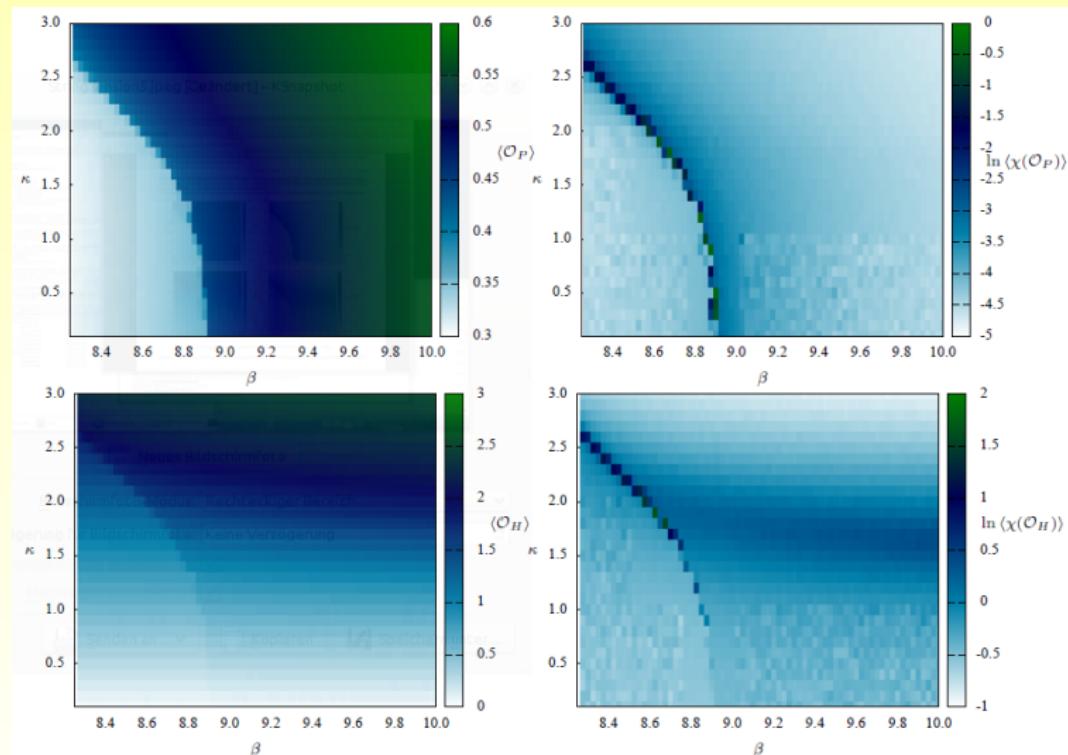


FIG. 4: Average plaquette, Higgs action and susceptibilities near the critical point on  $6^3 \times 2$  lattice.



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# Results

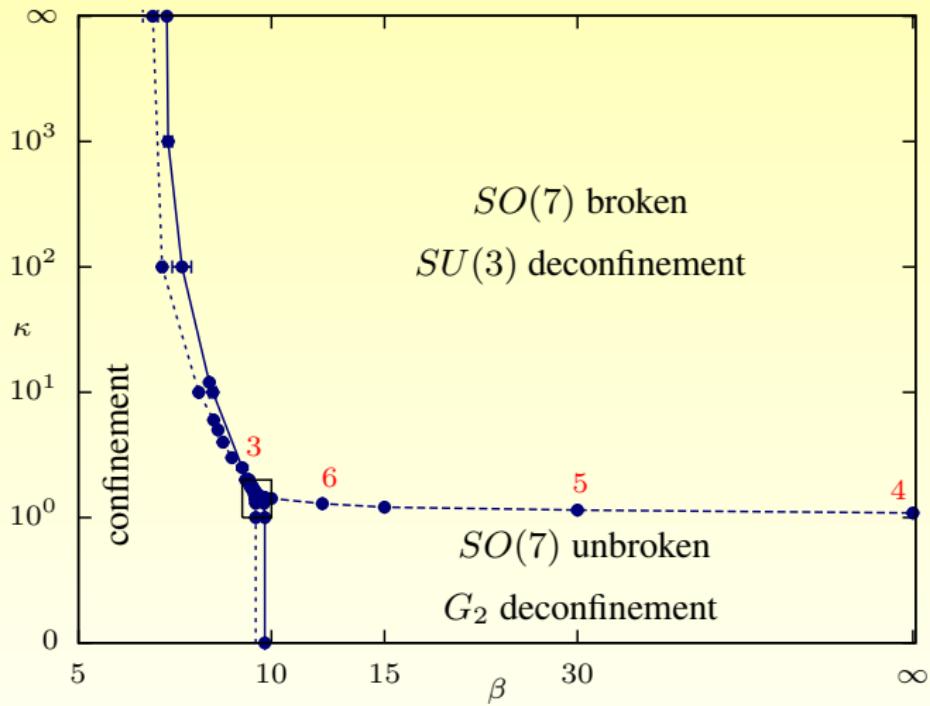
- average plaquette action, Higgs action and Polyakov loop
- susceptibilities (and higher derivatives) and finite size analysis
- large  $\beta \propto 1/g^2 \Rightarrow$  Higgs transition line
  - cluster algorithm for SO(7) nonlinear sigma model
  - line of second order PT  $O(7) \rightarrow O(6)$
- line of first order PT  $G_2 \rightarrow SU(3)$  with small gap in between
- triple point

$$\beta_{\text{crit}} = 9.55(5) \quad , \quad \kappa_{\text{crit}} = 1.50(4)$$

- first order (almost?) line hits second order line
- confining phase meets two deconfining phases



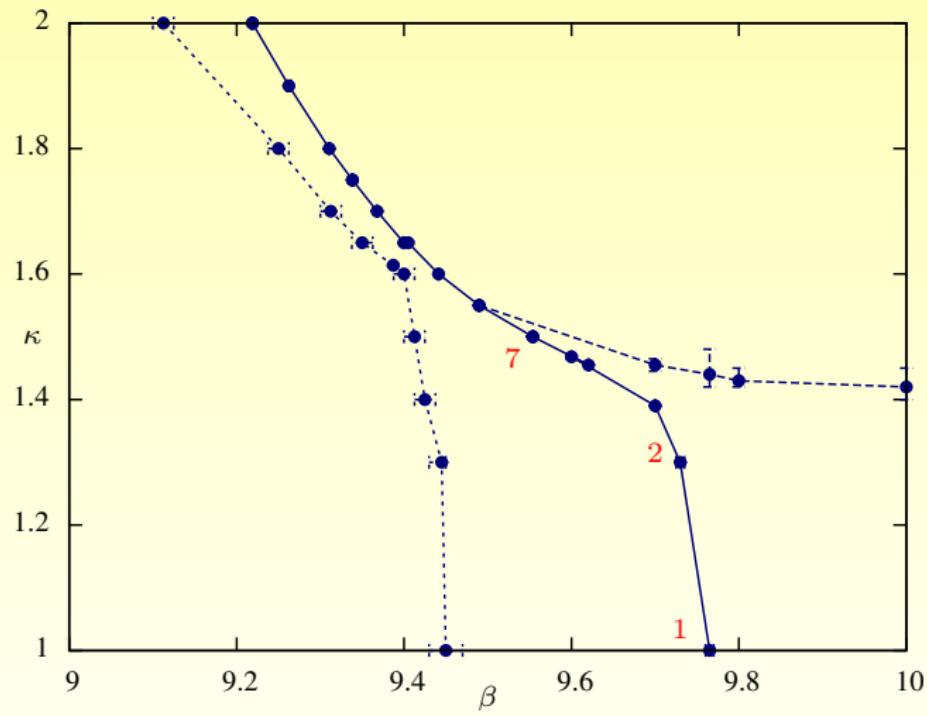
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phase diagram ( $16^3 \times 6$ ): global picture



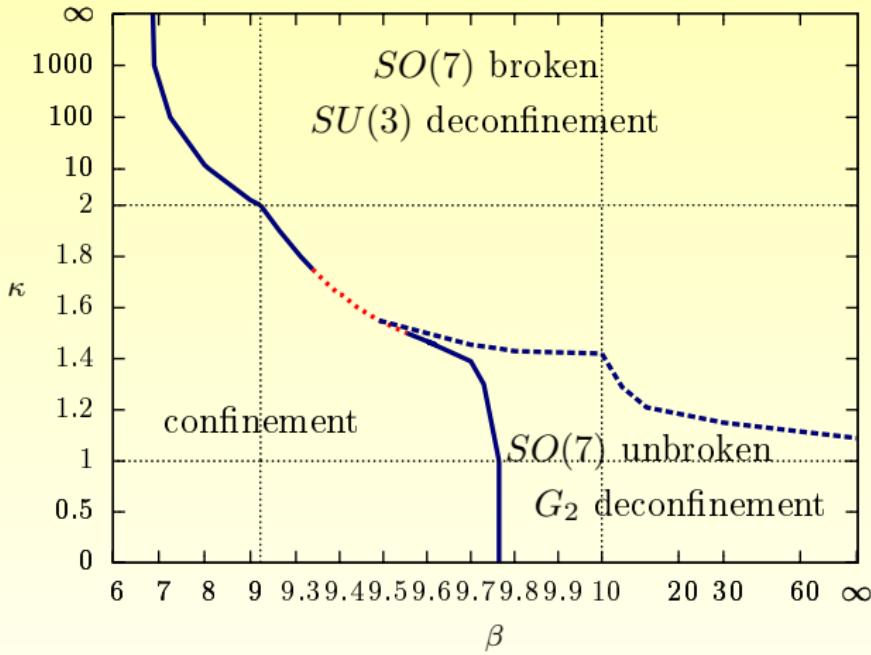
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phase diagram ( $16^3 \times 6$ ): where the lines almost meet



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## phases of $G_2$ YMH-theory

Wellegehausen, Wozar, AW



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# $G_2$ QCD with dynamical fermions

collaboration with

Axel Maas, Lorenz von Smekal und Bjoern Wellegehhausen

- fermionic determinant real and positive
- no sign problem: simulations at finite  $T$  and  $\mu$
- expected particle spectrum:  
glueballs  
bosonic quark-quark bound states (mesons, diquarks)  
fermionic 3 quark states (baryons)  
fermionic 1 quark - 2 gluon bound states (fermion hybrids)



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- confinement-deconfinement transition (Polyakov loop)
- chiral symmetry breaking (chiral condensate)
- quenched: same critical temperatures
- $G_2$  has fermionic baryons

Maas, Gatringer

degenerate Fermigas at large  $\rho_B$

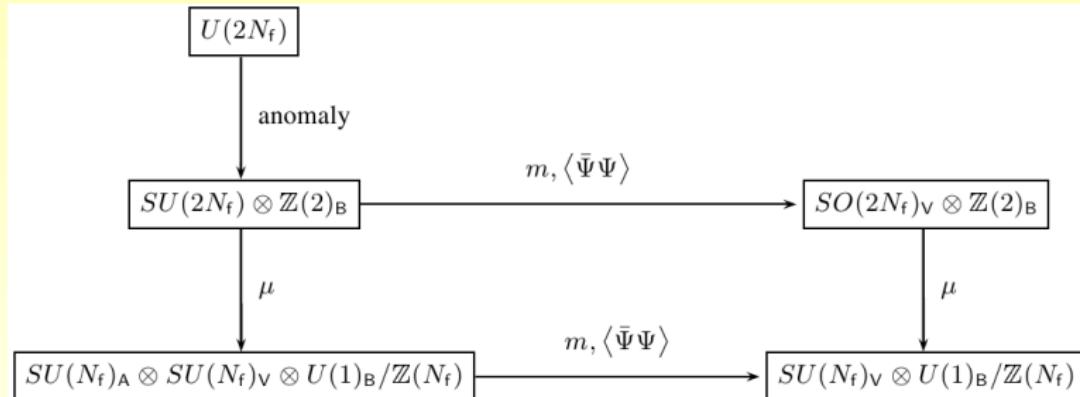
relevant for physics of compact 'stellar objects'

- Bose-condensates of diquarks, ...
- results from (expensive) numerical simulations



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# $G_2$ -QCD: global symmetries and breaking

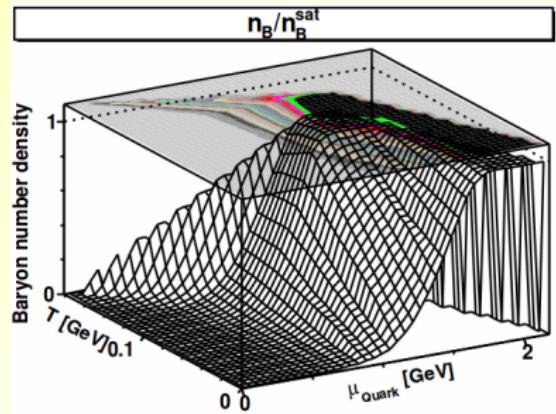
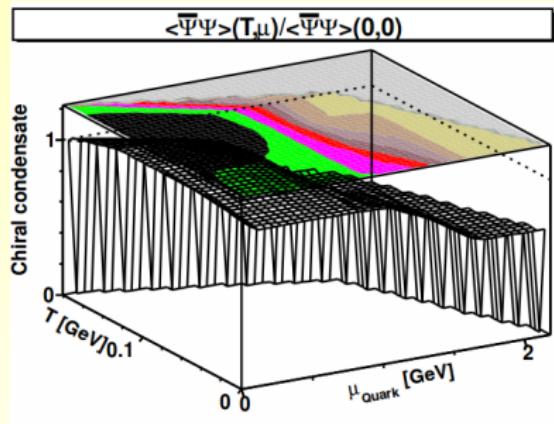


- anomalous
- spontaneous  $\langle \bar{\psi} \psi \rangle$
- explicit  $m, \mu$



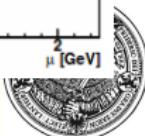
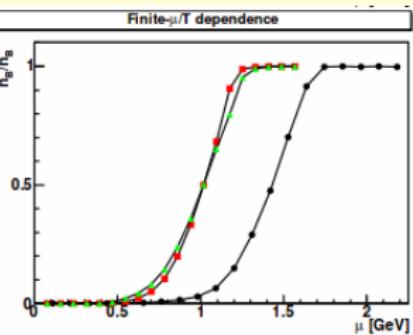
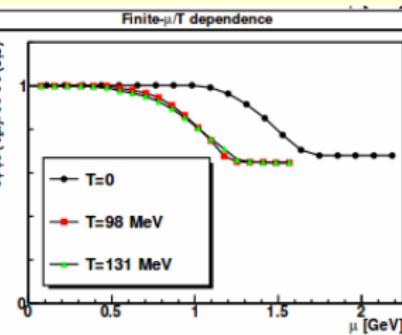
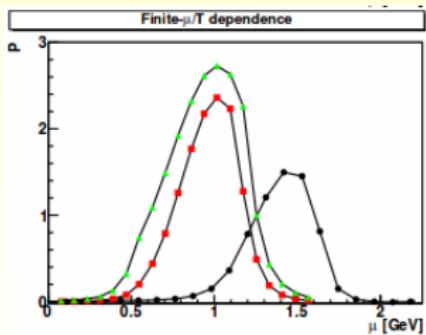
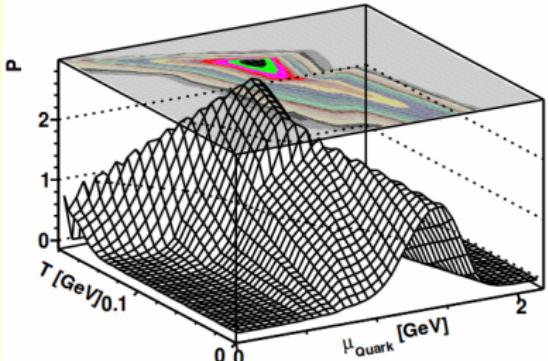
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# first results on $\langle\bar{\psi}\psi\rangle$ , $n_B$ and $\langle P \rangle$



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## Polyakov loop



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# Spectroscopy I

mesons (baryon number 0)

| Name     | $\mathcal{O}$                  | $T$  | J | P | C |
|----------|--------------------------------|------|---|---|---|
| $\pi$    | $\bar{u}\gamma_5 d$            | SASS | 0 | - | + |
| $\eta$   | $\bar{u}\gamma_5 u$            | SASS | 0 | - | + |
| $a$      | $\bar{u}d$                     | SASS | 0 | + | + |
| $f$      | $\bar{u}u$                     | SASS | 0 | + | + |
| $\rho$   | $\bar{u}\gamma_\mu d$          | SSSA | 1 | - | + |
| $\omega$ | $\bar{u}\gamma_\mu u$          | SSSA | 1 | - | + |
| $b$      | $\bar{u}\gamma_5 \gamma_\mu d$ | SSSA | 1 | + | + |
| $h$      | $\bar{u}\gamma_5 \gamma_\mu u$ | SSSA | 1 | + | + |

exotic particles (baryon number 1)

| Name      | $\mathcal{O}$  | $T$  | J   | P     | C     |
|-----------|--|------|-----|-------|-------|
| $N'$      | $T^{abc}(\bar{u}_a\gamma_5 d_b)u_c$                                      | SAAA | 1/2 | $\pm$ | $\pm$ |
| $\Delta'$ | $T^{abc}(\bar{u}_a\gamma_\mu u_b)u_c$                                    | SSAS | 3/2 | $\pm$ | $\pm$ |
| Hybrid    | $\epsilon_{abcdefg} u^a F_{\mu\nu}^{bc} F_{\mu\nu}^{de} F_{\mu\nu}^{fg}$ | SSSS | 1/2 | $\pm$ | $\pm$ |

$$T: (x, s, C, F)$$



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## diquarks (baryon number 2)

| Name        | $\mathcal{O}$  | $T$  | J | P | C |
|-------------|--|------|---|---|---|
| $d(0^{++})$ | $\bar{u}^c \gamma_5 u + c.c.$  | SASS | 0 | + | + |
| $d(0^{+-})$ | $\bar{u}^c \gamma_5 u - c.c.$  | SASS | 0 | + | - |
| $d(0^{-+})$ | $\bar{u}^c u + c.c.$   | SASS | 0 | - | + |
| $d(0^{--})$ | $\bar{u}^c u - c.c.$   | SASS | 0 | - | - |
| $d(1^{++})$ | $\bar{u}^c \gamma_\mu d - \bar{d}^c \gamma_\mu u + c.c.$                   | SSSA | 1 | + | + |
| $d(1^{+-})$ | $\bar{u}^c \gamma_\mu d - \bar{d}^c \gamma_\mu u - c.c.$                   | SSSA | 1 | + | - |
| $d(1^{-+})$ | $\bar{u}^c \gamma_5 \gamma_\mu d - \bar{d}^c \gamma_5 \gamma_\mu u + c.c.$ | SSSA | 1 | - | + |
| $d(1^{--})$ | $\bar{u}^c \gamma_5 \gamma_\mu d - \bar{d}^c \gamma_5 \gamma_\mu u - c.c.$ | SSSA | 1 | - | - |

## baryons (baryon number 3)

| Name     | $\mathcal{O}$                              | $T$  | J   | P     | C     |
|----------|--|------|-----|-------|-------|
| $N$      | $T^{abc} (\bar{u}_a^c \gamma_5 d_b) u_c$   | SAAA | 1/2 | $\pm$ | $\pm$ |
| $\Delta$ | $T^{abc} (\bar{u}_a^c \gamma_\mu u_b) u_c$ | SSAS | 3/2 | $\pm$ | $\pm$ |



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# spectroscopy II

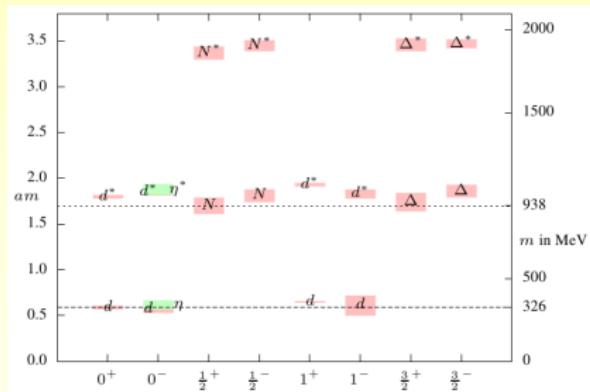
- diquark masses are degenerate
- contain only connected contributions (as pions in QCD)
- $m_\eta - m_{\text{diquark}} = \text{disconnected contributions}$
- tree level improved Szymanzik action
- Wilson fermions (chiral properties?)
- no sources for diquarks needed
- $N_F$  complex-valued pseudo-fermions plus RHMC
- two time-scale integration (Sexton-Weingarten and leapfrog)
- further optimizations (preconditioning, adaptive mesh, ...)



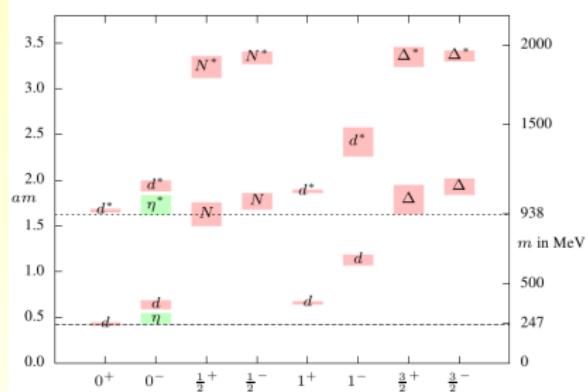
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# masses of mesons, diquarks, baryons

| Ensemble | $\beta$ | $\kappa$ | $m_{d(0+)} a$ | $m_N a$  | $m_{d(0+)} [\text{MeV}]$ | $a [\text{fm}]$ | $a^{-1} [\text{MeV}]$ | MC |
|----------|---------|----------|---------------|----------|--------------------------|-----------------|-----------------------|----|
| Heavy    | 1.05    | 0.147    | 0.59(2)       | 1.70(9)  | 326                      | 0.357(33)       | 552(50)               | 7K |
| Light    | 0.96    | 0.159    | 0.43(2)       | 1.63(13) | 247                      | 0.343(45)       | 575(75)               | 5K |



heavy ensemble

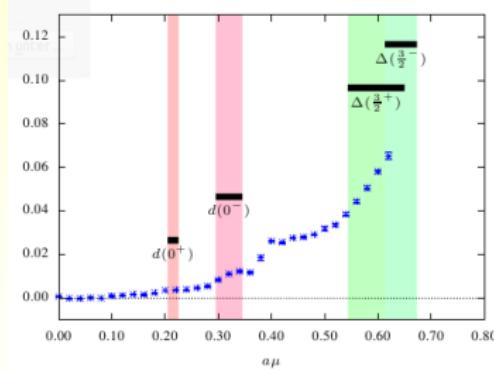
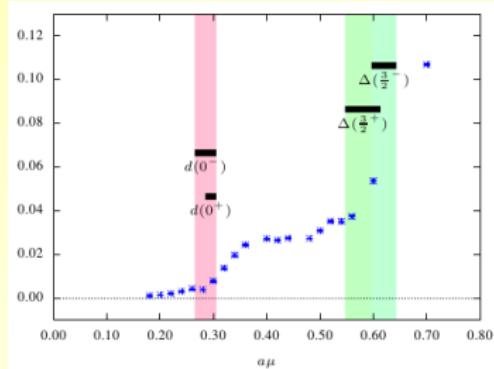


light ensemble

Wellegehausen, Maas, Smekal, AW (2013)



# zooming in in: baryon density vs. chemical potential



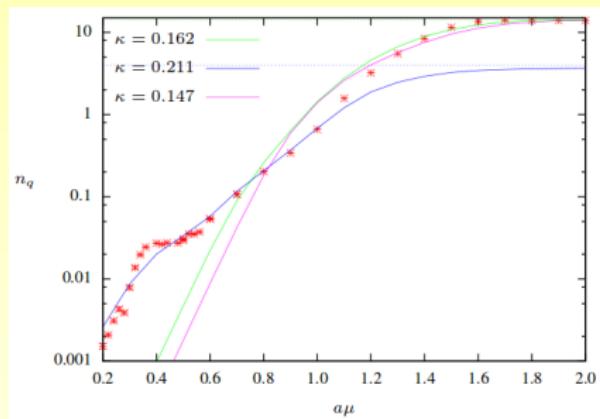
- $n_q$  grows rapidly at half of  $(0^-)$  mass (silver blaze)
- plateaus visible for larger  $n_q$
- three transitions
- phase between  $\mu_q = 300 - 600$  MeV:  
hadronic phase  
→ quasi particle picture

Welleghausen, Maas, Smekal, AW (2013)



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# comparison with fermi gas of (Wilson) fermions



- fit above  $a\mu = 1 \Rightarrow \kappa = 0.162, n_f^{\text{sat}} = 14.4$   
cp. with free quarks:  $\kappa = 0.147, n_f^{\text{sat}} = 14 \Rightarrow$  saturation regime
- fit below  $a\mu = 1 \Rightarrow \kappa = 0.211, n_f^{\text{sat}} = 4.02$   
cp. lattice gas of freee  $\Delta$ -baryons:  $n_f^{\text{sat}} = 4$
- $0.6 \leq a\mu \leq 1$ :  $n_q$  due to fermionic baryons (same as spectroscopy)



## (preliminary) interpretation

- low density: in accordance with **silver blaze**  
clean signal (no diquark sources)
- two small jumps at **diquark thresholds**  
⇒ two (probably) second order PT?
- two **plateaus** after thresholds
- **Bose-condensates** of diquarks?  
admixture with gas of diquarks?
- one (probably) first order PT at  $\approx \Delta$  threshold  
simulations slow down near PT
- hadronic phase for higher  $n_q$  (under investigation)
- $a\mu \gtrsim 1$ : lattice artifacts, e.g. saturation effects



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# Summary

- $G_2$  QCD is a useful laboratory
- accessible at finite density by lattice methods
- phases and transition at high densities and temperatures
- condensates, access to hadronic phase
- access to dense baryonic matter (as in  $n^*$ )
- shares many features with real QCD
- interpolation  $G_2$ -QCD  $\rightarrow$  QCD with Higgs-mechanism possible
- full phase diagram in principle accessible to simulations



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# Outlook

- clarify further nature of dense and cold phases
- include finite temperature effects beyond rough overview
- follow first order transition line (critical end point?)
- break  $G_2$ -QCD  $\rightarrow$  QCD with quarks via Higgs-mechanism
- deformation vs. sign problem?
- testbed for model building
- testbed for alternative approaches (eg. renormalization group)?

Thanks for your attention



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