# Exercise 9: Perturbative RG and the renormalized potential

In this exercise we want to give a functional representation based on a renormalized potential to our lecture's results based on perturbation theory for small coupling.

Recall that during the lecture we promoted the standard  $\phi^4$  Lagrangian in d = 4 dimensions

$$\mathcal{L}_{4} = \frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi + \frac{1}{2}g_{2}\phi^{2} + \frac{1}{4!}g_{4}\phi^{4}$$

to a d-dimensional action by introducing a scale  $\mu$  as

$$\mathcal{L}_d = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi + \frac{1}{2}g_2\phi^2 + \frac{1}{4!}\mu^{4-d}g_4\phi^4$$

The "promotion" was done in such a way that the field  $\phi$  has always canonical dimension, but the couplings  $g_2$  and  $g_4$  have maintained the same dimension when  $4 \to d$ . The potential is for now restricted to be a polynomial of the fourth order.

## Part 1:

Consider the generalization of  $\mathcal{L}_4$  to a full functional potential

$$\mathcal{L}_4 = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + V(\phi)$$

The original Lagrangian can be recovered by expanding  $V(\phi) = \frac{1}{2}g_2\phi^2 + \frac{1}{4!}g_4\phi^4$  and the couplings are the Taylor coefficients of this expansion. Find the generalization of the potential that is needed to promote the Lagrangian to  $\mathcal{L}_d$  and leaves the dimension of the couplings invariant.

#### Hint:

You may try the ansatz

$$\mathcal{L}_d = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \mu^A V(\mu^B \phi)$$

and determine the unknown constants A and B by comparing with the coupling's Lagrangian.

After 1) dimensional regularization, 2) minimal subtraction of the  $\frac{1}{\epsilon}$  poles and 3) an unimportant rescaling of factors of  $4\pi$ , the beta function of the coupling  $g_4$  at two loops in d = 4 is

$$\beta_{g_4} = 3g_4^2 - \frac{17}{3}g_4^3$$

and the anomalous dimension is

$$\eta = \frac{1}{6}g_4^2$$

#### Part 2:

Assume that in the functional representation the potential has beta function

$$\beta_V = C_1 \left( V''(\phi) \right)^2 + C_2 V''(\phi) \left( V'''(\phi) \right)^2 + \frac{1}{2} \eta \phi V'(\phi)$$

and the anomalous dimension  $\eta$  has the form

$$\eta = C_3 \left( V^{(4)}(0) \right)^2$$

Determine the constants  $C_1$ ,  $C_2$  and  $C_3$ .

## Hint:

There is no need to compute them from diagrams if you substitute  $V(\phi) = \frac{1}{4!}g_4\phi^4 + \ldots$  and interpret  $\beta_V$  as the generator of the beta functions  $\beta_V = \frac{1}{4!}\beta_{g_4}\phi^4 + \ldots$ . It is probably convenient to determine the constant  $C_3$  of the anomalous dimension first.

Now we are interested in computing the critical exponents *below* four dimensions using the beta functions and the perturbative expansion of d = 4.

One trick to do this is to "trade" the expansion in  $g_4$  for a  $\varepsilon = 4 - d$  expansion. We use the fact that  $V(\phi)$  and  $\phi$  are canonically normalized in d dimensions and switch to a dimensionless potential  $v(\varphi) = \mu^{-d}V(\mu^{d/2-1}\varphi)$ . (Notice that the field  $\phi$  is already renormalized and the contribution of  $\eta$  is already inside  $\beta_V$  above so there is no need to rescale by a further wavefunction.)

We are purposedly distinguishing between  $d = 4 - \varepsilon$  (which parametrizes any dimension below the upper critical one for statistical field theory) and  $d = 4 - \epsilon$  (which analytically continued the theory to make it finite). After the new additional rescaling all couplings are dimensionless and the beta function is

$$\beta_{v} = -dv(\varphi) + \frac{d-2}{2}\varphi v'(\varphi) + \beta_{V}|_{V \to v}$$

$$= -dv(\varphi) + \frac{d-2+\eta}{2}\varphi v'(\varphi) + \frac{1}{2}\left(v''(\varphi)\right)^{2} - \frac{1}{2}v''(\phi)\left(v'''(\phi)\right)^{2}$$

$$\eta = -\frac{1}{6}(v^{(4)}(0))^{2}$$

(This expression has a familiar scaling term that should remind you of another exercise. It should also suggest you the correct values of  $C_1$ ,  $C_2$  and  $C_3$  as a check of part 2!) The fixed point solutions of  $\beta_v = 0$  are of the form  $v^*(\varphi) = \frac{1}{4!}\lambda_4^*(\varepsilon)\varphi^4$ .

# (Optional) Part 3:

- Find all the fixed points  $\lambda_4^*(\varepsilon)$  up to order  $\varepsilon^2$ . How many are they and why?
- Which fixed point is the one that we need for the  $\varepsilon$  expansion and why?
- Find a simple way to determine the critical exponent  $\eta$  and  $\nu$  to order  $\varepsilon^2$ . (The results should be  $\eta = \frac{\varepsilon^2}{54}$  and  $\nu = \frac{1}{2} + \frac{\varepsilon}{12} + \frac{7\varepsilon^2}{162}$ .)