

Exercise 4: The effective action

Consider a free massive scalar field in d dimensions with action

$$S[\varphi] = \frac{1}{2} \int d^d x \{(\partial\varphi)^2 + m^2\varphi^2\} = \frac{1}{2} \int d^d x \varphi (-\partial^2 + m^2) \varphi$$

and the path integral

$$e^{-\Gamma[\varphi]} = \int D\chi e^{-S[\varphi+\chi]}$$

The physical interpretation of the path integral above goes as follows: the field χ are **fluctuations** over a **background field** φ . By integrating the fluctuations we are obtaining an effective action for the field configuration φ which is valid at $\langle\chi\rangle = 0$.

Part 1:

Use the path-integral to show that the effective action is

$$\Gamma[\varphi] = S[\varphi] + \frac{1}{2} \text{Tr} \log \mathcal{O}$$

in which $\mathcal{O} = -\partial^2 + m^2$.

Hints for part 1:

The action is free and therefore the dependences on χ and φ are separable (exactly like when we did momentum shell RG in class). Use the Gaussian integral formula for the functional determinant.

Part 2:

Apply the Laplace transform method to show

$$\frac{1}{2} \text{Tr} \log \mathcal{O} = \frac{1}{2} \int_0^\infty ds \mathcal{L}^{-1}[f](s) \text{Tr} e^{-s\mathcal{O}}$$

in which $\mathcal{L}^{-1}[f](s)$ is the inverse Laplace transform of the function $f(x) = \log x$.

Hints for part 2:

Given a function $f(x)$, the relation with its inverse Laplace transform is

$$f(x) = \int_0^\infty ds \mathcal{L}^{-1}[f](s) e^{-sx}$$

Part 3:

Show that the inverse Laplace transform of the logarithm is

$$\mathcal{L}^{-1}[\log](s) = -\frac{1}{s}$$

and therefore

$$\frac{1}{2} \text{Tr} \log \mathcal{O} = -\frac{1}{2} \int_0^\infty \frac{ds}{s} \text{Tr} e^{-s\mathcal{O}}$$

Hints for part 3:

The relation of the inverse Laplace transform with the original function is

$$\mathcal{L}^{-1}[f](s) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} dz f(z) e^{sz}$$

with $\gamma \in \mathbb{R}$ such that all the poles z_i of $f(z)$ lie on the left part of the complex plane with respect to γ : $\gamma > \text{Re}(z_i) \forall i$. You might want to use the following property of the inverse Laplace transform:

$$\mathcal{L}^{-1}[f'](s) = -s\mathcal{L}^{-1}[f](s),$$

with f' the derivative of f .

Part 4:

Use the momentum space representation of the operator

$$e^{-s\mathcal{O}} |p\rangle = e^{-s(p^2+m^2)} |p\rangle$$

for a normalized state $|p\rangle$, and of the trace

$$\text{Tr}(\dots) = \int \frac{d^d p}{(2\pi)^d} \langle p | (\dots) | p \rangle$$

to find an explicit formula for

$$\frac{1}{2} \text{Tr} \log \mathcal{O}$$

(integrate first in p and then in s). What are the assumptions that you have to make on d and m^2 for the integral to be convergent? When do you have to make these assumptions? What happens if $m^2 = 0$?

Hints for part 4:

Recall the integral form of the Euler Gamma function:

$$\Gamma(z+1) = \int_0^\infty dt e^{-t} t^{z-1},$$

which is convergent for $z > 0$.
