

Phases of generalized Potts-Models and their Relevance for Gauge Theories

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O'Raifeartaigh Symposium

Budapest

22.-24. June 2006

Potts-Models

Polyakov-Loop Dynamics

Gluodynamics and Potts-Models

Modified mean field approximation

Results of MC-simulations

Conclusions

Potts-Models

Polyakov-Loop
Dynamics

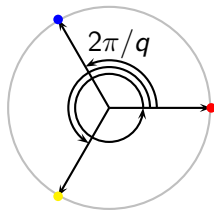
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generalized Ising models:

$$\theta_x \in \{2\pi k/q\}, \quad 1 \leq k \leq q$$

$$H = -J \sum_{\langle xy \rangle} \cos(\theta_x - \theta_y)$$

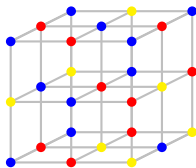
$$\mathbb{Z}_q : \theta_x \rightarrow \theta_x + 2\pi n/q$$

- ▶ ferromagnetic phase: q ground states

phase transition symmetric \leftrightarrow ferromagnetic

$d = 2$: second order $q \leq 4$, first order $q > 4$

$d = 3$: second order $q \leq 2$, first order $q > 2$



anti-ferromagnetic phase:

rich vacuum structures

symmetric \leftrightarrow antiferrom:

$d = 3, q = 3$: second order

entropy of ground states?

- entropy $S_B(p) = - \sum p(w) \log p(w) \Rightarrow$ free energy

$$\beta F = \inf_p (\beta \langle H \rangle_p - S_B) \Rightarrow p_{\text{Gibbs}} \sim e^{-\beta H}$$

- variational characterization of (convex) effective action:

$$\Gamma[m] = \inf_p \left(\beta \langle H \rangle_p - S(p) \mid \langle e^{i\theta(x)} \rangle_p = m(x) \right)$$

- mean field approximation:

$$\rho(w) = \prod_x \rho_x(\theta_x) \Rightarrow \Gamma_{\text{MF}}[m]$$

translational invariance: $p_x = p \Rightarrow m(x) = m$

effective potential: $\Gamma_{\text{MF}}[m] = V u_{\text{MF}}(m)$

$$u_{\text{MF}}(m) = \inf_p \left(-Kmm^* + \sum_{\theta} p(\theta) \log p(\theta) \right)$$

$$m = \sum_{\theta} p(\theta) e^{i\theta}, \quad K = dJ.$$

▶ antiferromagnetic phase:

translational invariance on sublattices $\Lambda = \Lambda_1 \cup \Lambda_2$

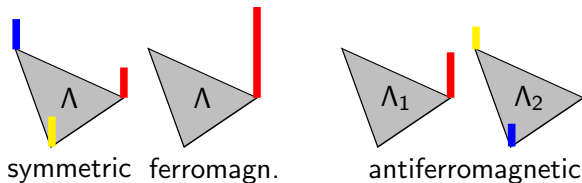
two neighbours in different sublattices

$p(x) = p_i \Rightarrow m(x) = m_i$ for $x \in \Lambda_i$

$$u_{\text{MF}}(m_1, m_2) = \frac{1}{2} \left(K |m_1 - m_2|^2 + \sum_i u_{\text{MF}}(m_i) \right),$$

▶ $K > K_{f,c} > 0 \Rightarrow m_1 = m_2 \neq 0$, \mathbb{Z}_q -broken

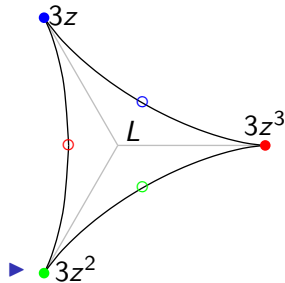
$K < K_{a,c} < 0 \Rightarrow m_1 \neq m_2 \neq 0$, \mathbb{Z}_{2q} -broken



Polyakov-Loop Dynamics

- ▶ finite temperature gluodynamics
order parameter for confinement: Polyakov loop
effective action:

$$e^{-S_{\text{eff}}[\mathcal{P}]} = \int \mathcal{D}U \delta \left(\mathcal{P}_{\mathbf{x}}, \prod_{t=0}^{N_t} U_{t,\mathbf{x};0} \right) e^{-S_w[U]}$$



gauge invariance:

$$S_{\text{eff}} = S_{\text{eff}}[L], \quad L_{\mathbf{x}} = \text{Tr} \mathcal{P}_{\mathbf{x}}$$

global Z_3 center symmetry:

$$S_{\text{eff}}[L] = S_{\text{eff}}[z \cdot L]$$

good ansatz for S_{eff} ?

- ▶ **strong coupling expansion** for $S_{\text{eff}}[\mathcal{P}]$
 $\Rightarrow \mathbb{Z}_3$ -invariant character expansion
nearest neighbour interaction

$$S_{\text{eff}} = \lambda_{10} S_{10} + \lambda_{21} S_{21} + \lambda_{20} S_{20} + \lambda_{11} S_{11} + \dots$$

$$S_{10} = \sum (\chi_{10}(\mathcal{P}_x) \chi_{01}(\mathcal{P}_y) + h.c.), \quad S_{21} = \dots$$

- ▶ **center-transformation:**

$$\chi_{pq}(z\mathcal{P}) = z^{p-q} \chi_{pq}(\mathcal{P}), \quad z^3 = z^* z = 1$$

With $L = \text{Tr} \mathcal{P}$: leading terms

$$\begin{aligned} S_{\text{eff}} &= (\lambda_{10} - \lambda_{21}) \sum (L_x L_y^* + h.c.) \\ &+ \lambda_{21} \sum (L_x^2 L_y + L_y^2 L_x + h.c.) \end{aligned}$$

- ▶ complex field with **compact target space**, \prod (reduced Haar measures), close relation to 3-state Potts model

Gluodynamics and Potts-Models

- ▶ naive reduction to Potts: $\mathcal{P}_x \rightarrow e^{i\theta_x} \mathbb{1} \in \text{centre}$

$$S_{\text{eff}} \rightarrow H \quad \text{with} \quad J = 18(\lambda_{01} + 4\lambda_{21})$$

true for all $S_{\text{eff}} \Rightarrow S_{\text{eff}}$ is extension of \mathbb{Z}_3 model.

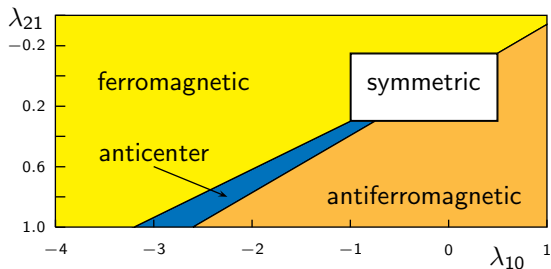
- ▶ Conjecture (Svetitsky, Yaffe):

effective finite-temperature $SU(N)$ -gluodynamics in
 d dimensions $\cong \mathbb{Z}_N$ spin model in $d - 1$ dimensions.

- ▶ same critical exponents $SU(2)$ and Ising (Engels et.al)
same universality class (symmetric \leftrightarrow ferrom.)

	β/ν	γ/ν	ν
4d $SU(2)$	0.545	1.93	0.65
3d Ising	0.516	1.965	0.63

- ▶ relevance for finite temperature $SU(N)$ with $N > 2$?
transition first order! → phase diagrams
- ▶ classical analysis: minimize S_{eff}



- ▶ quantum fluctuations \Rightarrow include symmetric phase
new ferromagnetic **anti-center phase**
qualitatively correct phase diagram

Modified mean field approximation

- ▶ variational characterisation of Γ :
fix $\langle \chi_j(\mathcal{P}_x) \rangle$ for all χ_j in S_{eff}
- ▶ **mean field** approximation \Rightarrow product measure

$$\mathcal{DP} \longrightarrow \prod_x d\mu_{\text{red}}(\mathcal{P}_x) \rho_x(\mathcal{P}_x)$$

- ▶ translational invariance on sublattices in $\Lambda = \Lambda_1 \cup \Lambda_2$
 \Rightarrow nontrivial variational problem on two-sites
- ▶ most simple effective model (Polonyi)

$$S_{\text{eff}} = \lambda S_{10} = \lambda \sum (L_x L_y^* + \text{h.c.})$$

Lagrangean multiplier for \bar{L}_i on Λ_i

- ▶ mean field effective potential for minimal model

$$2u_{\text{MF}}(L_1, L_1^*, L_2, L_2^*) = -d\lambda|L_1 - L_2|^2 + \sum v_{\text{MF}}(L_i, L_i^*)$$

$$v_{\text{MF}}(L, L^*) = d\lambda|L|^2 + \gamma_0(L, L^*)$$

γ_0 Legendre-transform of

$$w_0(j, j^*) = \log \int d\mu_{\text{red}} \exp(jL + j^* L^*)$$

- ▶ order parameters:

$$L = \frac{1}{2}(L_1 + L_2), \quad M = \frac{1}{2}(L_1 - L_2), \quad \ell = |L|, \quad m = |M|.$$

- ▶ group integral in **closed form** not known for $SU(3)$!!
 $\int \exp(j \text{Tr}(U)) = \text{hypergeometric function}$

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Potts-Models

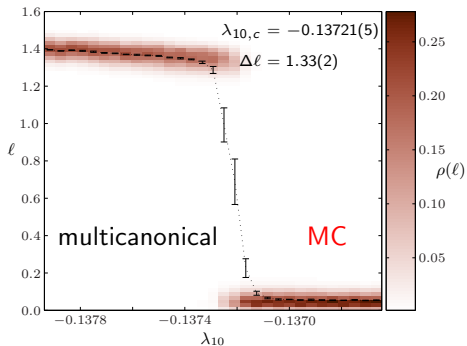
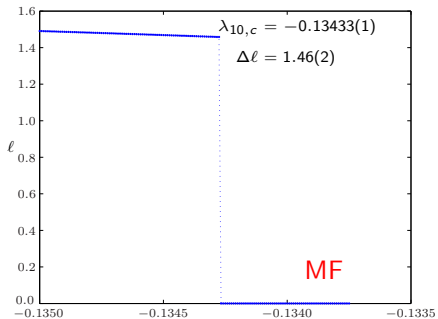
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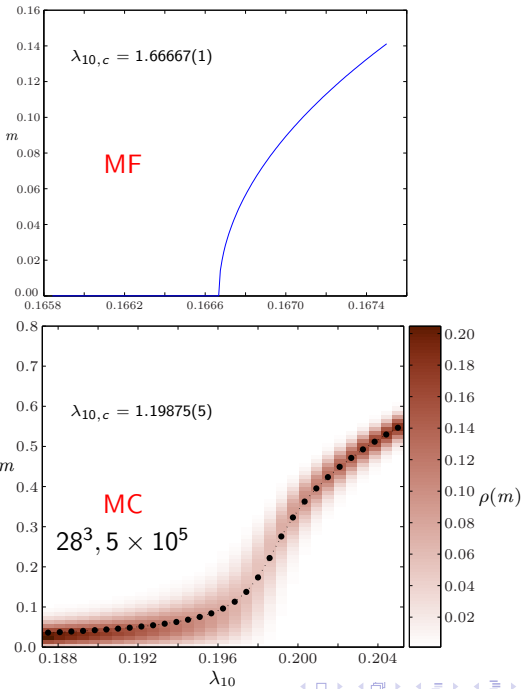
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- ▶ Why is mean field so good?
conjecture: 3 = upper crit. dimension for 3-state potts
- ▶ critical exponents of $S \leftrightarrow AF$:

exponent	3-state Potts	minimal S_{eff}
ν	0.664(4)	0.68(2)
γ/ν	1.973(9)	1.96(2)

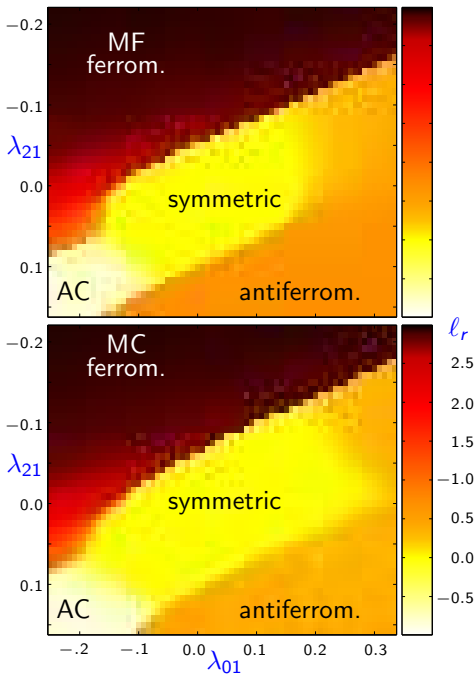
critical exponents in mean field?

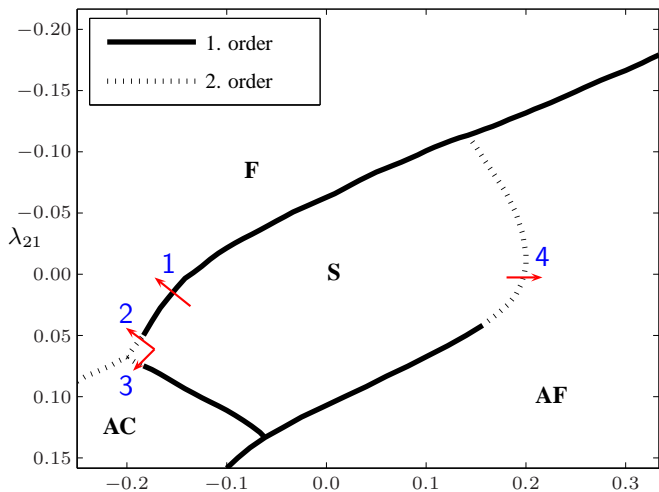
- ▶ finite temperature gluodynamics
 - effective \mathbb{Z}_3 models with compact target spaces
 - 3-state Potts-model

universality test in 'unphysical region'
(for gluodynamics)

Results of MC-Simulations

- ▶ phase diagram and transitions → histograms
large statistics, expensive → fast algorithms!
standard Metropolis: 5% to 10% accuracy
- ▶ multicanonical algorithm: up to 20^3 lattices near first order transitions
- ▶ new cluster algorithm near second order transitions:
auto-correlation times down by two orders of magnitude on larger lattices
- ▶ comparison with mean field results for two-coupling (costly).
- ▶ rich phase structure: 4 different phases, second und first order transitions, tricritical points(?), mean field very good.





jenLaTT, Linux cluster, 8000 MC simulations, 3000 CPUh

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Potts-Models

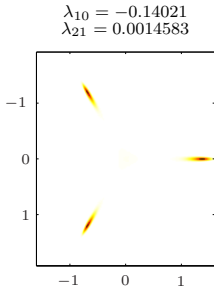
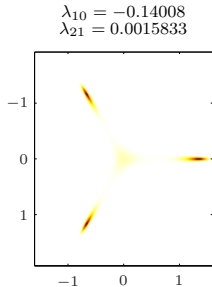
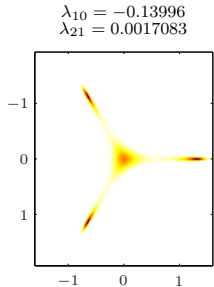
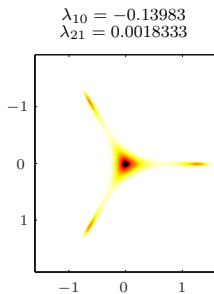
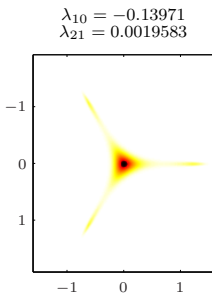
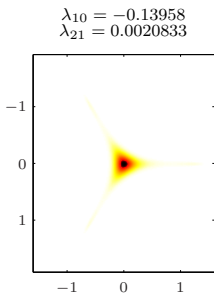
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• 1: Histogramm of $L, S \leftrightarrow$ FM, 1st

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Potts-Models

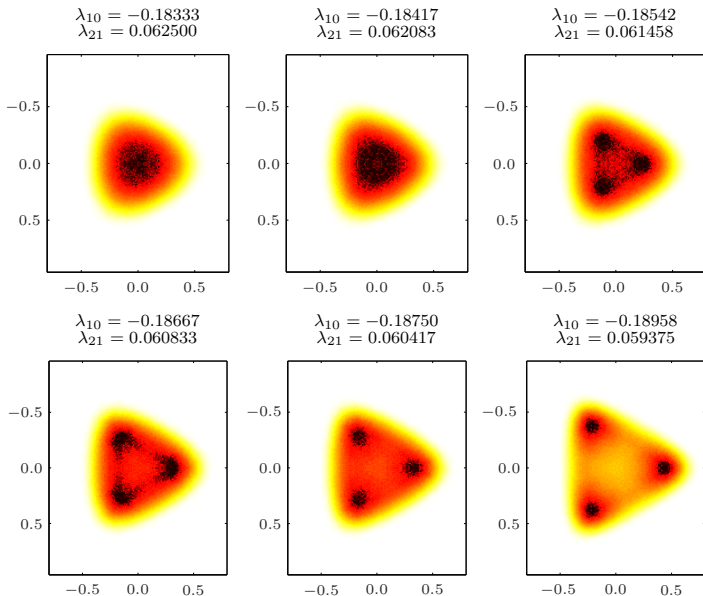
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• 2: Histogramm of $L, S \leftrightarrow FM$, 2nd

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Potts-Models

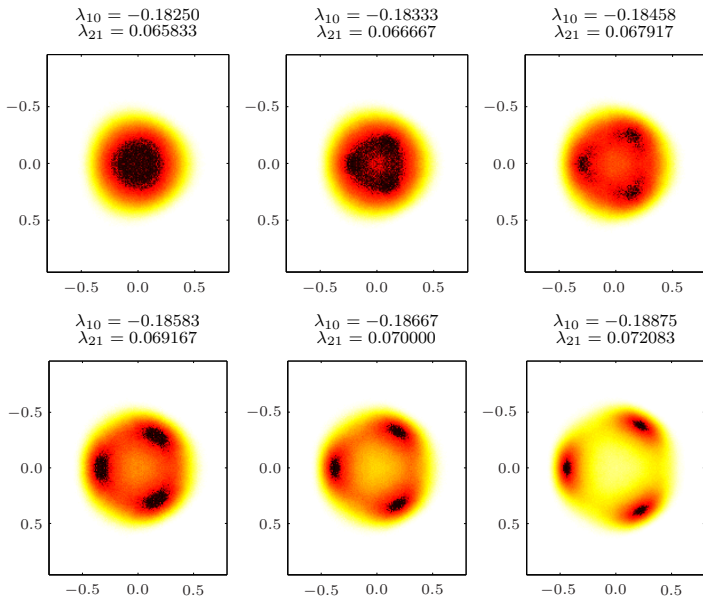
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• 3: Histogramm of $L, S \leftrightarrow AC$, 2nd

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Potts-Models

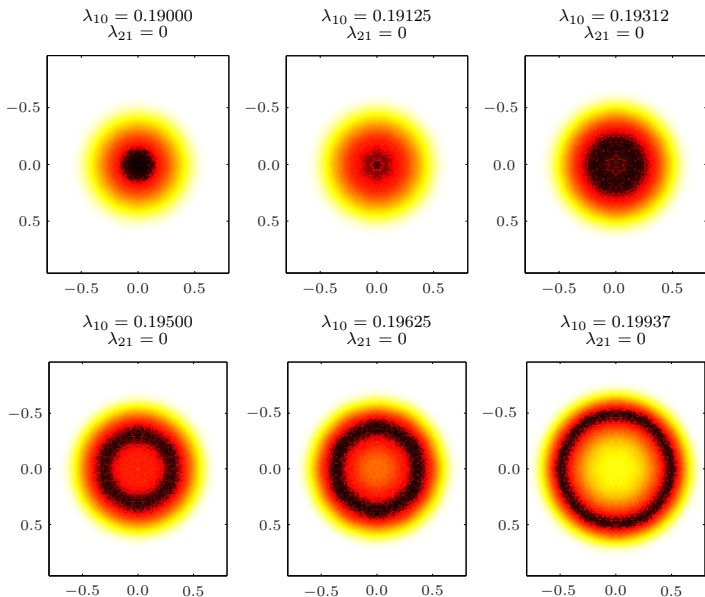
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•4: Histogramm of $M, S \leftrightarrow$ AF, 2nd

Conclusions

- ▶ strong coupling for Polyakov-loops effective action
 - ▶ modified MF for non-translationally invariant states
 - ▶ new efficient cluster algorithm!
-
- ▶ rich phase structure for simple \mathbb{Z}_3 -models
 - ▶ mean field unexpectedly accurate ($d_c = 3?$)
-
- ▶ calculate $\lambda_j(\beta)$ via IMC (cp. $SU(2)$)
 - ▶ efficient 'group-theoretic' Schwinger-Dyson equations!
 - ▶ vacuum-sector of AF phase? $SU(N)$ group integrals?
 - ▶ is AC-phase relevant for gluodynamics?
 - ▶ include fermions in effective Polyakov-loop dynamics.

JHEP 06 (2004) 005, PRD 72 (2005) 065005, hep-lat/0605012